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Inverse Problem for a Nonlinear Pseudiparabolic Differential Equation with Final Condition and Gerasimov–Caputo Operator

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Abstract. *In this paper in rectangle domain an inverse problem for a fractional analogue of the pseudoparabolic differential operator with mixed conditions, degeneration and identification source is considered. Fractional operator is the Gerasimov–Caputo type and the solution of the nonlinear differential equation with two spatial variables is studied in the class of generalized functions. The nonlinear Fourier series method is used and by the aid of Kilbas–Saigo function a nonlinear countable system of functional integral equation is obtained. In the proof of unique solvability of the countable system is applied the method of successive approximations in combination with the method of compressing mapping. We use the Cauchy–Schwarz inequality and the Bessel inequality in proving the absolute and uniform convergence of the obtained Fourier series. Then we derive the desire redefinition function also in the form of Fourier series.*

Key words: *Inverse problem, fractional analogue, nonlinear pseudoparabolic differential equation, final condition, degeneration, identification source.*

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Introduction. Problem statement

Some problems of mechanics turn out to be initial-boundary (mixed) problems. Many mixed problems are studied in solving different problems of hydrodynamics [1]. In [2], [3] mixed problems for linear differential equations of parabolic and hyperbolic types were studied. In works [4, 5], mixed problems for nonlinear differential and integro-differential equations of the

second, fourth and higher orders were studied. The main equations of the theory of non-stationary filtration in fractured-pore formations are formulated in the work of G. I. Barenblatt, Yu. P. Zheltov and I. N. Kochina [6] (see also [7]) and, further, developed by many authors [8]–[13]. The theory and applications of fractional calculus have been developed by many authors (see, for example, [14]–[22]). Investigation of the well-known fractional order differentiation operators of Riemann–Liouville type and Gerasimov–Caputo type are important, because they describe diffusion processes [15, vol. 1, 47–85]. A physical and engineering interpretations of the generalized fractional operators are given in [15, vol. 4–8]. Note that boundary value conditions of the type (0.3) (see, below) take place in modeling problems of the flow around a profile by a subsonic velocity stream with a supersonic zone. Different mixed and boundary value problems for differential and integro-differential equations with identification source were studied in the works of many authors (see, for example, in the works [23]–[43]).

In the domain $\Omega = \{(t, x, y) | 0 < t < T, 0 < x, y < l\}$ a partial differential equation of the following form

$$\left[{}_C D_{0t}^\alpha - {}_C D_{0t}^\alpha \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) - t^\beta \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \right] U(t, x, y) = f(t, x, y, U(t, x, y)) \quad (0.1)$$

is considered with final point integral condition

$$U(T, x, y) = \varphi(x, y) + \int_0^T R(s)U(s, x, y)ds, \quad 0 \leq x, y \leq l, \quad (0.2)$$

where β , T and l are given positive real numbers, for $0 < \alpha \leq 1$ [15, Vol. 1, p. 34]

$${}_C D_{0t}^\alpha \eta(t) = I_{0t}^{1-\alpha} \eta'(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t \frac{\eta'(s)}{(t-s)^\alpha} ds, \quad {}_C D_{0t}^1 \eta(t) = \eta'(t), \quad t \in (0, T)$$

is Gerasimov–Caputo type fractional operator [44, p. 112]

$$I_{0t}^\alpha \eta(t) = \frac{1}{\Gamma(\alpha)} \int_0^t \frac{\eta(s)ds}{(t-s)^{1-\alpha}}, \quad t \in (0, T)$$

is Riemann–Liouville integral operator, $a(t), b(t) \in C[0, T]$ is given function, $f(x, y)$ is given function, $\varphi(x, y) \in L_2([0, l]^2)$ is redefinition function. Sometimes, we use the notations: $\Omega_T \equiv [0, T]$, $\Omega_l \equiv [0, l]$.

We denote by $W_{2,t,x}^{n,4nm}(\bar{\Omega})$ the class of continuously functions $U(t, x)$ of two variables in closed rectangle $\bar{\Omega} = \{(t, x) | 0 \leq t \leq T, 0 \leq x \leq l\}$ and having partial derivatives in it

$$\frac{\partial U(t, x)}{\partial x}, \frac{\partial^2 U(t, x)}{\partial x^2}, \dots, \frac{\partial^{4nm} U(t, x)}{\partial x^{4nm}}, \frac{\partial U(t, x)}{\partial t}, \dots, \frac{\partial^n U(t, x)}{\partial t^n},$$

each of them belongs $L_2(\bar{\Omega})$, and n, m are positive integers.

Problem 0.1. Find the generalized solution $U(t, x, y) \in W_{2,t,x,y}^{2,\alpha,2,2}(\bar{\Omega})$, which satisfies partial differential equation (0.1), final value integral condition (0.2), boundary value conditions

$$U(t, 0, y) = U(t, l, y) = U(t, x, 0) = U(t, x, l) = 0 \quad (0.3)$$

almost everywhere.

Problem 0.2. Find the pair of functions $U(t, x, y) \in W_{2,t,x,y}^{2,\alpha,2,2}(\bar{\Omega})$ $\varphi(x, y) \in L_2([0, l]^2)$, first of which satisfies the additional condition

$$\int_0^T H(s)U(s, x, y)ds = \psi(x, y), \quad (0.4)$$

where $\psi(x, y) \in L_2([0, l]^2)$ is given function and $\psi(0, y) = \psi(l, y) = \psi(x, 0) = \psi(x, l) = 0$. We assume that for the given functions are true the following boundary conditions

$$\varphi(0, y) = \varphi(l, y) = \varphi(x, 0) = \varphi(x, l) = 0,$$

$$f(t, 0, y, U(t, 0, y)) = f(t, l, y, U(t, l, y)) = f(t, x, 0, U(t, x, 0)) = f(t, x, l, U(t, x, l)) = 0.$$

1 Reducing the problem 0.1 into countable system of nonlinear integral equations

Nontrivial solutions of the problem (0.1)–(0.3) are sought as a Fourier series

$$U(t, x, y) = \sum_{n,m=1}^{\infty} u_{n,m}(t) \vartheta_{n,m}(x, y) \quad (1.1)$$

and we suppose that

$$f(t, x, y, U) = \sum_{n,m=1}^{\infty} f_{n,m}(t, \cdot) \vartheta_{n,m}(x, y), \quad (1.2)$$

where

$$u_{n,m}(t) = \int_0^l \int_0^l U(t, \eta, \xi) \vartheta_{n,m}(\eta, \xi) d\eta d\xi, \quad f_{n,m}(t, \cdot) = \int_0^l \int_0^l f(t, \eta, \xi, \cdot) \vartheta_{n,m}(\eta, \xi) d\eta d\xi, \quad (1.3)$$

$$\vartheta_{n,m}(x, y) = \frac{2}{l} \sin \lambda_n x \sin \lambda_m y = \frac{2}{l} \sin \frac{\pi n}{l} x \sin \frac{\pi m}{l} y, \quad n, m = 1, 2, \dots$$

Substituting Fourier series (1.1) and (1.2) formally into partial differential equation (0.1), we obtain a countable system of ordinary fractional differential equations of $0 < \alpha < 1$ -order with degeneration

$${}_C D_{0t}^\alpha u_{n,m}(t) + \mu_{n,n}^2 t^\beta u_{n,m}(t) = \frac{1}{1 + \lambda_{n,m}^2} f_{n,m} \left(t, \sum_{i,j=1}^{\infty} u_{i,j}(t) \vartheta_{i,j}(x, y) \right), \quad (1.4)$$

where

$$\mu_{n,m}^2 = \frac{\lambda_{n,m}^2}{1 + \lambda_{n,m}^2}, \quad \lambda_{n,m}^2 = \left(\frac{\pi n}{l} \right)^2 (n^2 + m^2).$$

We suppose also that there is true the following Fourier expansion

$$\varphi(x, y) = \sum_{n,m=1}^{\infty} \varphi_{n,m} \vartheta_{n,m}(x, y),$$

where

$$\varphi_{n,m} = \int_0^l \int_0^l \varphi(\eta, \xi) \vartheta_{n,m}(\eta, \xi) d\eta d\xi. \quad (1.5)$$

By Fourier coefficients (1.3) and (1.5), the final value integral condition (0.2) takes the form

$$u_{n,m}(T) = \varphi_{n,m} + \int_0^T R(s)u_{n,m}(s)ds. \quad (1.6)$$

We use the well known Kilbas–Saigo type function, which is generalized two-parametric Mittag–Leffler function [15, Vol. 1. pp. 269–295]

$$E_{\alpha,\beta}(z) = \sum_{m=0}^{\infty} \frac{z^m}{\Gamma(\alpha m + \beta)}, \quad z, \alpha, \beta \in \mathbb{C}, \quad \operatorname{Re}(\alpha) > 0.$$

The Kilbas–Saigo function is defined for real $\alpha, m, l \in \mathbb{R}$ and complex $l \in \mathbb{C}$ is defined by the following form [14, 16]

$$E_{\alpha,m,l}(z) = \sum_{k=0}^{\infty} c_k z^k, \quad c_0 < 1, \quad c_k = \prod_{j=1}^{k-1} \frac{\Gamma(\alpha[jm + l] + 1)}{\Gamma(\alpha[jm + l + 1] + 1)}, \quad k = 1, 2, \dots$$

The Kilbas–Saigo functions belong to the class of entire functions in the complex plane [45].

Let us consider the final value problem for a countable system of ordinary differential equation of fractional order with degeneration

$${}_C D_{0t}^{\alpha} u_{n,m}(t) = -\mu_{n,m}^2 t^{\beta} u_{n,m}(t) + f_{n,m}(t, \cdot), \quad u_{n,m}(0) = h_{n,m}, \quad (1.7)$$

where $\beta \in \mathbb{R}$, $0 < \mu_{n,m}^2 < 1$,

$$h_{n,m} = \int_0^l \int_0^l h(\eta, \xi) \vartheta_{n,m}(\eta, \xi) d\eta d\xi, \quad h(x, y) \in L_2([0, l]^2),$$

$$f_{n,m}(t, \cdot) = \int_0^l \int_0^l f \left(t, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}(t) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi.$$

Let $\gamma \in [0, 1)$. Then we consider the class of following functions

$$C_{\gamma}[0, T] = \{g_{n,m}(t) : t^{\gamma} g_{n,m}(t) \in C[0, T]\},$$

$$C_{\gamma}^{\alpha}[0, T] = \{g_{n,m}(t) \in C[0, T] : {}_C D_{0t}^{\alpha} g_{n,m}(t) \in C_{\gamma}[0, T]\}.$$

We use the known fact that for $|\arg z| \leq \sigma$ and $|z| \geq 0$ the following estimate is true [45]

$$|E_{\alpha,\mu}(z)| \leq M_1(1 + |z|)^{\frac{1-\delta}{\alpha}} e^{\operatorname{Re} z \frac{1}{\alpha}} + \frac{M_2}{1 + |z|}, \quad (1.8)$$

where M_1 and M_2 are constants, not depending from z ; $\alpha < 2$, $z \in \mathbb{C}$, δ is real constant and σ is fixed number from the interval $\left(\frac{\pi\alpha}{2}, \min\{\pi, \pi\alpha\}\right)$. If we put $z = \mu_{n,m}^2 t^\beta (t - \tau)^\alpha$, $\delta = \alpha$, then from (1.8) we have

$$\begin{aligned} \left| E_{\alpha, \alpha} \left(\mu_{n,m}^2 t^\beta (t - \tau)^\alpha \right) \right| &\leq M_1 \left[1 + \mu_{n,m}^2 t^\beta (t - \tau)^\alpha \right]^{\frac{1-\alpha}{\alpha}} e^{(\mu_{n,m}^2 t^\beta (t - \tau)^\alpha)^{\frac{1}{\alpha}}} + \\ &+ \frac{M_2}{1 + \mu_{n,m}^2 t^\beta (t - \tau)^\alpha} \leq M_3. \end{aligned} \quad (1.9)$$

Lemma 1.1. *Let $\gamma \in [0, \alpha]$, $\beta \geq 0$. Then for all $g_{n,m}(t) \in C_\gamma[0, T]$ there exists a unique solution $u_{n,m}(t) \in C_\gamma^\alpha[0, T]$ of the Cauchy problem (1.7). This solution has the following form*

$$u_{n,m}(t) = \varphi_{n,m} E_{\alpha, 1 + \frac{\beta}{\alpha}, \frac{\beta}{\alpha}} \left(-\mu_{n,m}^2 t^{\alpha+\beta} \right) + \int_0^t K(t, \tau) g_{n,m}(\tau) d\tau, \quad (1.10)$$

where

$$K(t, \tau) = \sum_{i=1}^n K_i(t, \tau), \quad (1.11)$$

$$K_0(t, \tau) = \frac{1}{\Gamma(\alpha)} (t - \tau)^{\alpha-1}, \quad K_i(t, \tau) = \frac{\mu_{n,m}^2}{\Gamma(\alpha)} \int_\tau^t s^\beta (t - s)^{\alpha-1} K_{i-1}(s, \tau) ds, \quad i = 1, 2, \dots, \quad (1.12)$$

$E_{\alpha, 1 + \frac{\beta}{\alpha}, \frac{\beta}{\alpha}} \left(\mu_{n,m}^2 t^{\alpha+\beta} \right)$ is Kilbas–Saigo function.

Moreover, for the kernel (1.11) in the case of $\gamma \in [0, \alpha]$, $\beta \geq 0$ there holds the following estimate

$$|K(t, \tau)| \leq (t - \tau)^{\alpha-1} E_{\alpha, \alpha} \left(\mu_{n,m}^2 t^\beta (t - \tau)^\alpha \right) \leq (t - \tau)^{\alpha-1} M_3, \quad (1.13)$$

where $M_3 = \text{const}$.

Consequently, the general solution of the countable system of ordinary differential equations (1.7) we write as

$$u_{n,m}(t) = C_{n,m} E_{\alpha, 1 + \frac{\beta}{\alpha}, \frac{\beta}{\alpha}} \left(-\mu_{n,m}^2 t^{\alpha+\beta} \right) + \frac{1}{1 + \lambda_{n,m}^2} \int_0^t K(t, s) f_{n,m}(s, \cdot) ds, \quad (1.14)$$

where $C_{n,m}$ is arbitrary constant, kernel $K(t, s)$ is defined by the formulas (1.11), (1.12) and for the kernel is true the estimate (1.13).

To find the unknown coefficients $C_{n,m}$ in (1.14), we use condition (1.6). So, substituting the equation (1.14) into (1.6), we obtain

$$\begin{aligned} u_{n,m}(T) &= C_{n,m} E_{\alpha, 1 + \frac{\beta}{\alpha}, \frac{\beta}{\alpha}} \left(-\mu_{n,m}^2 T^{\alpha+\beta} \right) + \frac{1}{1 + \lambda_{n,m}^2} \int_0^T K(T, s) f_{n,m}(s, \cdot) ds = \\ &= \varphi_{n,m} + \omega \int_0^T R(s) u_{n,m}(s) ds. \end{aligned}$$

Hence, we derive

$$C_{n,m} = \varphi_{n,m} Q_T^{-1} + Q_T^{-1} \int_0^T \left\{ R(t) u_{n,m}(t) - \frac{1}{1 + \lambda_{n,m}^2} K(T, t) f_{n,m}(t, \cdot) \right\} dt, \quad (1.15)$$

where

$$Q_T = E_{\alpha, 1 + \frac{\beta}{\alpha}, \frac{\beta}{\alpha}} \left(-\mu_{n,m}^2 T^{\alpha + \beta} \right). \quad (1.16)$$

Further, substituting the coefficient presentation's (1.15) into equation (1.14) and taking into account (1.16), we derive a countable system of nonlinear integral equations (CSNIE)

$$\begin{aligned} u_{n,m}(t) = & \varphi_{n,m} Q_t Q_T^{-1} + Q_t Q_T^{-1} \int_0^T R(s) u_{n,m}(s) ds + \\ & + \frac{1}{1 + \lambda_{n,m}^2} \int_0^T \bar{K}(t, s) f_{n,m}(s, \cdot) ds, \end{aligned} \quad (1.17)$$

where

$$\bar{K}(t, s) = \begin{cases} K(t, s) - Q_t Q_T^{-1} K(T, s), & t \leq s \leq T, \\ K(t, s), & 0 \leq s \leq t. \end{cases}$$

2 Redefinition function. Realisation of the inverse Problem

Substituting the countable system of integral equations (1.17) into Fourier series (1.1), we obtain

$$\begin{aligned} U(t, x, y) = & \sum_{n,m=1}^{\infty} \vartheta_{n,m}(x, y) \left\{ \varphi_{n,m} Q_t Q_T^{-1} + Q_t Q_T^{-1} \int_0^T R(s) u_{n,m}(s) ds + \right. \\ & \left. + \frac{1}{1 + \lambda_{n,m}^2} \int_0^T \bar{K}(t, s) f_{n,m}(s, \cdot) ds \right\}. \end{aligned} \quad (2.1)$$

By virtue of (1.1), (1.2) and (1.5), the additional condition (0.4) we rewrite in the form

$$\psi_{n,m} = \int_0^T H(s) u_{n,m}(s) ds.$$

We apply the last condition to the (1.17)

$$\begin{aligned} \psi_{n,m} = & \int_0^T H(s) \left\{ \varphi_{n,m} Q_s Q_T^{-1} + Q_s Q_T^{-1} \int_0^T R(\theta) u_{n,m}(\theta) d\theta + \right. \\ & \left. + \frac{1}{1 + \lambda_{n,m}^2} \int_0^T \bar{K}(s, \theta) f_{n,m}(\theta, \cdot) d\theta \right\} ds. \end{aligned} \quad (2.2)$$

Assume that $q_0 = \int_0^T H(s) Q_s Q_T^{-1} ds \neq 0$. Then from (2.2) we have

$$\begin{aligned} \varphi_{n,m} = & q_0^{-1} \psi_{n,m} - \frac{q_0^{-1}}{1 + \lambda_{n,m}^2} \int_0^T H(s) \int_0^T \overline{K}(s, \theta) f_{n,m}(\theta, \cdot) d\theta ds - \\ & - q_0^{-1} \int_0^T H(s) Q_s Q_T^{-1} \int_0^T R(\theta) u_{n,m}(\theta) d\theta ds. \end{aligned} \quad (2.3)$$

Hence, we obtain the Fourier series for redefinition function

$$\begin{aligned} \varphi(x, y) = & \sum_{n,m=1}^{\infty} \vartheta_{n,m}(x, y) \left\{ q_0^{-1} \psi_{n,m} - \frac{q_0^{-1}}{1 + \lambda_{n,m}^2} \int_0^T H(s) \int_0^T \overline{K}(s, \theta) f_{n,m}(\theta, \cdot) d\theta ds - \right. \\ & \left. - q_0^{-1} \int_0^T H(s) Q_s Q_T^{-1} \int_0^T R(\theta) u_{n,m}(\theta) d\theta ds \right\}. \end{aligned} \quad (2.4)$$

Redefinition function in (2.4) consists unknown function $u_{n,m}(t)$. So, the function (2.3) we substitute into the equation (1.17) and Fourier series (2.1):

$$\begin{aligned} u_{n,m}(t) = & J(t; u_{n,m}(t)) \equiv q_0^{-1} Q_t Q_T^{-1} \psi_{n,m} + \\ & + \frac{1}{1 + \lambda_{n,m}^2} \left[\int_0^T \overline{K}(t, s) \int_0^l \int_0^l F \left(t, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}(t) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi ds - \right. \\ & \left. - q_0^{-1} Q_t Q_T^{-1} \int_0^T H(s) \int_0^T \overline{K}(s, \theta) \int_0^l \int_0^l F \left(t, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}(t) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi d\theta ds \right] + \\ & + Q_t Q_T^{-1} \left[\int_0^T R(s) u_{n,m}(s) ds - q_0^{-1} \int_0^T H(s) Q_s Q_T^{-1} \int_0^T R(\theta) u_{n,m}(\theta) d\theta ds \right], \end{aligned} \quad (2.5)$$

$$\begin{aligned} U(t, x, y) = & \sum_{n,m=1}^{\infty} \vartheta_{n,m}(x, y) \left\{ q_0^{-1} Q_t Q_T^{-1} \psi_{n,m} + \right. \\ & + \frac{1}{1 + \lambda_{n,m}^2} \left[\int_0^T \overline{K}(t, s) \int_0^l \int_0^l F \left(t, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}(t) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi ds - \right. \\ & \left. - q_0^{-1} Q_t Q_T^{-1} \int_0^T H(s) \int_0^T \overline{K}(s, \theta) \int_0^l \int_0^l F \left(\theta, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}(\theta) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi d\theta ds \right] + \\ & \left. + Q_t Q_T^{-1} \left[\int_0^T R(s) u_{n,m}(s) ds - q_0^{-1} \int_0^T H(s) Q_s Q_T^{-1} \int_0^T R(\theta) u_{n,m}(\theta) d\theta ds \right] \right\}. \end{aligned} \quad (2.6)$$

3 Main theorems

We will use the concepts of the following well-known Banach spaces. Space $B_2[0, T]$ of sequences of continuous functions $\{u_{n,m}(t)\}_{n,m=1}^{\infty}$ on the segment $[0, T]$ with norm

$$\|\vec{u}(t)\|_{B_2[0,T]} = \sqrt{\sum_{n,m=1}^{\infty} \left(\max_{t \in [0,T]} |u_{n,m}(t)| \right)^2} < \infty.$$

Hilbert coordinate space ℓ_2 of number sequences $\{\varphi_{n,m}\}_{n,m=1}^{\infty}$ with norm

$$\|\vec{\varphi}\|_{\ell_2} = \sqrt{\sum_{n,m=1}^{\infty} |\varphi_{n,m}|^2} < \infty.$$

The space $L_2(\Omega_l^2)$ of square-summable functions on the domain $\Omega_l^2 = \Omega_l \times \Omega_l$ with norm

$$\|\vartheta(x, y)\|_{L_2(\Omega_l^2)} = \sqrt{\int_0^l \int_0^l |\vartheta(x, y)|^2 dx dy} < \infty.$$

It is not difficult to prove that Bessel's inequality is valid:

$$\|\vec{\varphi}\|_{\ell_2} = \sqrt{\sum_{n,m=1}^{\infty} |\varphi_{n,m}|^2} \leq \|\varphi(x, y)\|_{L_2(\Omega_l^2)}. \quad (3.1)$$

Theorem 3.1. *Let be fulfilled the smoothness conditions and following conditions:*

- 1). $\sum_{n,m=1}^{\infty} |\psi_{n,m}| < \infty$;
 - 2). $\max \left\{ \int_0^T |R(s)| ds; \int_0^T |H(s)| ds \right\} < \infty$;
 - 3). $\|f(t, x, y, u)\|_{L_2(\Omega \times (-\infty, \infty))} \leq C_0 = \text{const}$;
 - 4). $|f(t, x, y, u_1) - f(t, x, y, u_2)| \leq C(x, y) |u_1 - u_2|$, $0 < C(x, y) \in L_2(\Omega_l^2)$;
 - 5). $\rho = \overline{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 < 1$, where
- $\overline{M}_1 = \sum_{n,m=1}^{\infty} \frac{1}{\lambda_{n,m}^2}$, Δ_1 and Δ_2 will be defined below from (3.6) and (3.7), respectively.

Then CSNIE (2.5) has a unique solution in the space $B_2[0, T]$.

Proof. In proof of the theorem, we use the method of successive approximations:

$$\begin{cases} u_{n,m}^0(t) = q_0^{-1} Q_t Q_T^{-1} \psi_{n,m}, \\ u_{n,m}^{k+1}(t) = J(t; u_{n,m}^k), \quad k = 0, 1, 2, 3, \dots, \quad t \in [0, T]. \end{cases} \quad (3.2)$$

Let us estimate the zero approximation $u_{n,m}^0(t)$. We use the known estimate [14]

$$\frac{1}{1 + \Gamma(1 - \alpha)\tau} \leq E_{\alpha, m, m-1}(-\tau) \leq \frac{1}{1 + \frac{\Gamma(1 + \alpha(m-1))}{\Gamma(1 + \alpha m)}\tau},$$

which is true for every $\alpha \in [0, 1]$, $m > 0$ and $\tau \geq 0$. We put $\tau = -\mu_n^2 t^{\alpha+\beta}$, $m = 1 + \frac{\beta}{\alpha}$, $\Gamma(1 + \alpha(m-1)) = \Gamma(1 + \beta)$, $\Gamma(1 + \alpha m) = \Gamma(1 + \alpha + \beta)$, then we obtain

$$E_{\alpha, 1 + \frac{\beta}{\alpha}, \frac{\beta}{\alpha}} \left(-\mu_{n,m}^2 t^{\alpha+\beta} \right) \leq \frac{1}{1 + \frac{\Gamma(1 + \beta)}{\Gamma(1 + \alpha + \beta)} \mu_{n,m}^2 t^{\alpha+\beta}} \leq 1. \quad (3.3)$$

From (3.3) it is obvious that $0 < |Q_t Q_T^{-1}| < 1$. So, taking into account (1.13), for the first approximation from (3.2) we have

$$\begin{aligned} \|\bar{u}^0(t)\|_{B_2[0,T]} &= \sqrt{\sum_{n,m=1}^{\infty} \left[\max_{0 \leq t \leq T} |u_{n,m}^0(t)| \right]^2} \leq \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} |u_{n,m}^0(t)| \leq \\ &\leq \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} |q_0^{-1}| |Q_t Q_T^{-1}| |\psi_{n,m}| \leq |q_0^{-1}| \sum_{n,m=1}^{\infty} |\psi_{n,m}| = \delta_0 < \infty. \end{aligned} \quad (3.4)$$

Due to the estimates (1.13), (3.4), applying the Cauchy–Shwartz inequality and Bessel's inequality (3.1), for the first difference we obtain

$$\begin{aligned} &\|\bar{u}^1(t) - \bar{u}^0(t)\|_{B_2[0,T]} \leq \\ &\leq \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} \frac{1}{\lambda_{n,m}^2} \left[\left| \int_0^T \bar{K}(t,s) \int_0^l \int_0^l F \left(t, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}^0(t) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi ds \right| + \right. \\ &+ \left. \left| q_0^{-1} \left| \int_0^T H(s) \int_0^T \bar{K}(s,\theta) \int_0^l \int_0^l F \left(t, \eta, \xi, \sum_{i,j=1}^{\infty} u_{i,j}^0(t) \vartheta_{i,j}(\eta, \xi) \right) \vartheta_{n,m}(\eta, \xi) d\eta d\xi d\theta ds \right| \right| + \right. \\ &+ \left. \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} \left[\left| \int_0^T R(s) |u_{n,m}^0(s)| ds \right| + \left| q_0^{-1} \left| \int_0^T H(s) Q_s Q_T^{-1} \int_0^T R(\theta) |u_{n,m}^0(\theta)| d\theta ds \right| \right] \right] \leq \\ &\leq M_1 \max_{0 \leq t \leq T} \|\bar{f}(t, \cdot)\|_{B_2[0,T]} \left\{ \int_0^T |\bar{K}(t,s)| ds + \left| q_0^{-1} \right| \int_0^T |H(s)| \int_0^T |\bar{K}(s,\theta)| d\theta ds \right\} + \\ &+ \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} |u_{n,m}^0(t)| \left\{ \int_0^T |R(s)| ds + \left| q_0^{-1} \right| \int_0^T |H(s)| \int_0^T |R(\theta)| d\theta ds \right\} \leq \\ &\leq \delta_0 [M_1 C_0 \Delta_1 + \Delta_2] = \delta_0 \delta_1 < \infty, \end{aligned} \quad (3.5)$$

where

$$\delta_1 = M_1 C_0 \Delta_1 + \Delta_2, \quad M_1 = \sqrt{\sum_{n,m=1}^{\infty} \frac{1}{\lambda_{n,m}^4}},$$

$$\Delta_1 = \int_0^T |\bar{K}(t, s)| ds + \left| q_0^{-1} \right| \int_0^T |H(s)| \int_0^T |\bar{K}(s, \theta)| d\theta ds, \quad (3.6)$$

$$\Delta_2 = \int_0^T |R(s)| ds + \left| q_0^{-1} \right| \int_0^T |H(s)| \int_0^T |R(\theta)| d\theta ds. \quad (3.7)$$

Due to the estimate (3.5), for the arbitrary difference $u_{n,m}^{k+1}(t) - u_{n,m}^k(t)$ we obtain the following estimate

$$\begin{aligned} & \left\| \bar{u}^{k+1}(t) - \bar{u}^k(t) \right\|_{B_2[0,T]} \leq \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} \left| u_{n,m}^{k+1}(t) - u_{n,m}^k(t) \right| \leq \\ & \leq \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} \frac{1}{\lambda_{n,m}^2} \left[\left| \int_0^T \bar{K}(t, s) \int_0^l \int_0^l C(\eta, \xi) \sum_{i,j=1}^{\infty} \left| u_{i,j}^k(s) - u_{i,j}^{k-1}(s) \right| \vartheta_{i,j}(\eta, \xi) \vartheta_{n,m}(\eta, \xi) d\eta d\xi ds \right| + \right. \\ & \left. + \left| q_0^{-1} \right| \left| \int_0^T H(s) \int_0^T \bar{K}(s, \theta) \int_0^l \int_0^l C(\eta, \xi) \sum_{i,j=1}^{\infty} \left| u_{i,j}^k(\theta) - u_{i,j}^{k-1}(\theta) \right| \vartheta_{i,j}(\eta, \xi) \vartheta_{n,m}(\eta, \xi) d\eta d\xi d\theta ds \right| \right] + \\ & + \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} \left[\left| \int_0^T |R(s)| \left| u_{i,j}^k(s) - u_{i,j}^{k-1}(s) \right| ds \right| + \left| q_0^{-1} \right| \left| \int_0^T H(s) \int_0^T |R(\theta)| \left| u_{i,j}^k(\theta) - u_{i,j}^{k-1}(\theta) \right| d\theta ds \right| \right] \leq \\ & \leq \frac{2}{l} \bar{M}_1 \max_{0 \leq t \leq T} \left\| C(x, y) \right\|_{B_2[0,T]} \left\| \bar{u}^k(t) - \bar{u}^{k-1}(t) \right\|_{B_2[0,T]} \Delta_1 + \sum_{n,m=1}^{\infty} \max_{0 \leq t \leq T} \left| u_{n,m}^k(t) - u_{n,m}^{k-1}(t) \right| \Delta_2 \leq \\ & \leq \delta_0 \delta_1 \left[\frac{2}{l} \bar{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 \right]^k = \delta_0 \delta_1 \rho^k, \end{aligned} \quad (3.8)$$

where

$$\rho = \frac{2}{l} \bar{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2, \quad \bar{M}_1 = \sum_{n,m=1}^{\infty} \frac{1}{\lambda_{n,m}^2}.$$

Since $\rho = \frac{2}{l} \bar{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 < 1$ and $\delta_0 < \infty$, $\delta_1 < \infty$, from estimates (3.8) follows that

$$\lim_{k \rightarrow \infty} \left\| \bar{u}^{k+1}(t) - \bar{u}^k(t) \right\|_{B_2[0,T]} = 0. \quad (3.9)$$

From estimates (3.4), (3.5) and (3.9) we obtain the existence and uniqueness of the solution $\bar{u}(t) \in B_2[0, T]$ to CSNIE (2.5). Theorem 3.1 is proved. \square

The Fourier series (2.6) we consider as a formal solution of the problem (0.1)–(0.3).

Theorem 3.2. *Let the conditions of Theorem 3.1 be satisfied and*

$$2\delta_0 \left[1 + \delta_1 \left(\frac{2}{l} \bar{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 \right) \right] < l.$$

If $\bar{u}(t) \in B_2[0, T]$ is the unique solution to CSNIE (2.5), then the series (2.6) will be generalized solution to the mixed problem (0.1)–(0.3) almost everywhere.

Proof. We consider the difference

$$\begin{aligned} \left| U(t, x, y) - U^k(t, x, y) \right| &= \left| \sum_{n,m=1}^{\infty} u_{n,m}^{\infty}(t) \vartheta_{n,m}(x, y) - \sum_{n,m=k}^{\infty} u_{n,m}^k(t) \vartheta_{n,m}(x, y) \right| \leq \\ &\leq \sum_{n,m=k}^{\infty} \left| u_{n,m}(t) \vartheta_{n,m}(x, y) \right|. \end{aligned}$$

As in the case of proof of the theorem 3.1, we obtain the estimate

$$\begin{aligned} \sum_{n,m=k}^{\infty} \left| u_{n,m}(t) \vartheta_{n,m}(x, y) \right| &\leq \sum_{n,m=k}^{\infty} \left| \vartheta_{n,m}(x, y) \right| \times \\ &\times \left\{ \left| q_0^{-1} \sum_{n,m=k}^{\infty} \left| \psi_{n,m} \right| + \delta_0 \delta_1 \left[\frac{2}{l} \overline{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 \right] \right\} \leq \\ &\leq \frac{2}{l} \delta_0 \left[1 + \delta_1 \left(\frac{2}{l} \overline{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 \right) \right] < \infty. \end{aligned} \quad (3.10)$$

Since $\vec{u}(t) \in B_2[0, T]$ and $\frac{2}{l} \delta_0 \left[1 + \delta_1 \left(\frac{2}{l} \overline{M}_1 \left\| C(x, y) \right\|_{L_2[0,l]^2} \Delta_1 + \Delta_2 \right) \right] < 1$ from (3.10) implies that

$$\lim_{k \rightarrow \infty} \left| U(t, x, y) - U^k(t, x, y) \right| = 0.$$

This proved the theorem 3.2. □

Now the function $u_{n,m}(t)$ in the representation (2.4) is known. As $\vec{u}(t) \in B_2[0, T]$, the proof of the convergence of series (2.4) is similarly to the proof of the Theorem 3.2.

Conclusion

In the domain $\Omega = \{(t, x, y) \mid 0 < t < T, 0 < x, y < l\}$ the generalized solvability of the inverse mixed problem (0.1)–(0.4) for a partial differential equation is considered for the case of α -order Gerasimov–Caputo type fractional operator with order $0 < \alpha \leq 1$. The considering equations depend from three independent arguments. First argument is time argument and with respect to this argument considering equation is fractional Gerasimov–Caputo type ordinary differential equation. Second and third arguments are spatial and the equations with respect to these arguments are differential equations of second order. The Fourier series method is used and a countable system of differential equations is obtained and studied. The generalized solvability of mixed problem (0.1)–(0.3) is derived in the form of the Fourier series.

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